



# Multiplicity associated to high $p_T$ events and percolation of color sources

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based on the work done in collaboration with J.Dias de Deus and C.Pajares

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## Multiplicity associated to high $p_T$ events and percolation of color sources

- What rare (self-shadowed) events are
- Universal relation among multiplicity distributions:  $P_C(n) = \frac{n}{\langle n \rangle} P(n)$
- Independent color sources scheme:  $\frac{\langle n \rangle_C}{\langle n \rangle} = 1 + \frac{D^2}{\langle n \rangle^2}$
- Introducing collectivity via percolation of color sources:  

$$\frac{\langle n \rangle_C}{\langle n \rangle} = 1 + \frac{1}{k}$$
- RESULT: clear bump for  $\frac{\langle n \rangle_C}{\langle n \rangle}$  at the centrality class for which the system is more disordered.

## Self-Shadowed events

- Considering a nucleon-nucleus collision as a superposition of independent nucleon-nucleon collisions:

$$\sigma^{hA}(b) = \sum_{n=1}^A \binom{A}{n} (\sigma T(b))^n (1 - \sigma T(b))^{A-n}$$

- Each independent elementary collision can only be of type C or not of type C in a binomial fashion:

$$\sigma = \sigma_C + \sigma_{NC}$$

$$(\sigma T(b))^n = \sum_{i=0}^n \binom{n}{i} (\sigma_C)^i (\sigma_{NC})^{n-i} T(b)^n$$

- At least one elementary interaction is of type C ( $C \otimes NC = C$ ):

$$\begin{aligned} \sigma_C^{hA}(b) &= \sum_{n=1}^A \binom{A}{n} \sum_{i=1}^n \binom{n}{i} (\sigma_C)^i (\sigma_{NC})^{n-i} T(b)^n \\ &= 1 - (1 - \sigma_C T(b))^A \end{aligned}$$

- $\sigma_C$  shadowed by itself,

$$\text{if } \sigma_C(b) \ll 1, \sigma_C^{hA} \simeq A\sigma_C \text{ if } \sigma_C(b) \simeq \sigma(b), \sigma_C^{hA} \simeq A^{2/3}$$

## Universal law for probability distributions

- Consider  $N(\nu)$  to be the total number of events produced with  $\nu$  collisions:

$$N(\nu) = \sum_{i=0}^{\nu} \binom{\nu}{i} (1 - \alpha_C)^{\nu-i} \alpha_C^i N(\nu)$$

- for small  $\alpha_C$  (rare event):

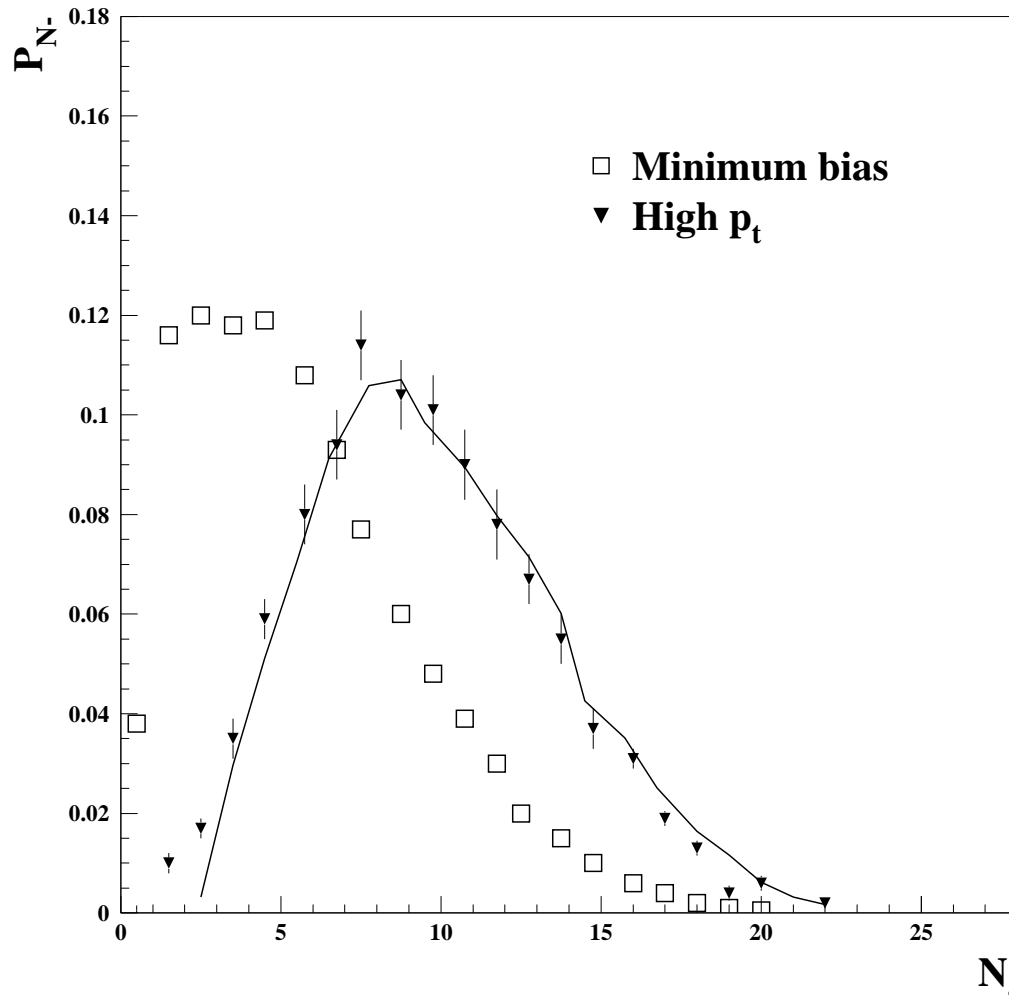
$$N_C(\nu) = \alpha_C \nu N(\nu) \quad \text{and} \quad N_{NC}(\nu) = (1 - \alpha_C \nu) N(\nu)$$

- For  $\sum_{\nu} N(\nu) = N$ ,  $\sum_{\nu} \nu N(\nu) = \langle \nu \rangle N$   
 $\sum_{\nu} N_C(\nu) = \sum_{\nu} \alpha_C \nu N(\nu)$ ,

- $P_C(\nu) = \frac{\alpha_C \nu N(\nu)}{\sum_{\nu} N_C(\nu)} = \frac{\nu N(\nu)}{\langle \nu \rangle \sum_{\nu} N(\nu)} = \frac{\nu}{\langle \nu \rangle} P(\nu)$  UNIVERSAL

## Universal law for probability distributions

- $P_C(\nu) = \frac{\nu}{\langle \nu \rangle} P(\nu)$
  
- KNO for particles is well approximated by KNO for collisions:
  - $\langle n \rangle P(n) = \langle \nu \rangle P(\nu)$
  
  - $n = \nu \bar{n}_0$
  - $\frac{D^2}{\langle n \rangle^2} = \frac{D(\nu)^2}{\langle \nu \rangle^2} + \frac{d^2}{\bar{n}_0^2 \langle \nu \rangle}$
  
  - assume dominance of the fluctuations in the number of collisions
  - the dominance increases with increasing  $\langle \nu \rangle$
  - the same kind of approximation for all other moments ...
  
- $P_C(n) = \frac{n}{\langle n \rangle} P(n) \longrightarrow$  also in  $E_T!$



- $P_C(n) = \frac{n}{\langle n \rangle} P(n)$
- Data for  $\alpha - \alpha$  at ISR
- C means  $p_T > 3\text{GeV}/c$

## Universal law for probability distributions

- $P_C(n) = \frac{n}{\langle n \rangle} P(n) \longrightarrow \frac{\langle n \rangle_C}{\langle n \rangle} = 1 + \frac{D^2}{\langle n \rangle^2}$

- dominance of the fluctuations in the number of collisions
- smallness of  $\alpha_C$
- linearity ( $N_C \propto \nu$ )

- and all based on the idea of superposition, independent emitters

and this is no longer valid at RHIC energies! for which collectivity manifests through the absence of scaling with collisions

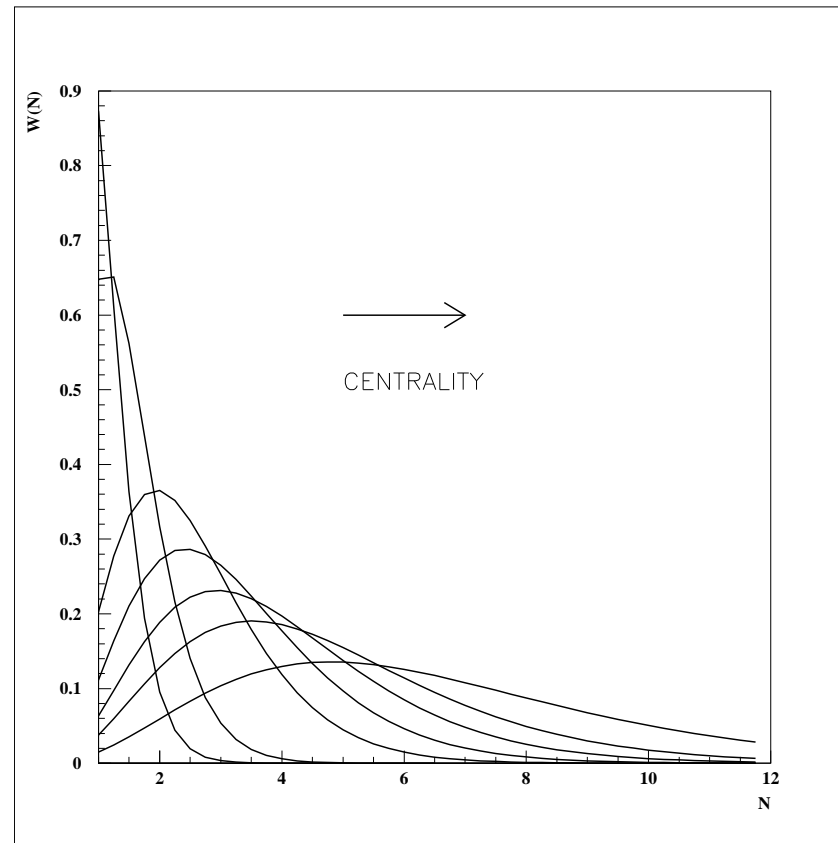
but eq still valid as we will see

## Percolation of color sources

- $P(n) = \int_0^\infty dN W(N) P(N, n)$
- $P(n)$ : particle distribution, the negative binomial is a good fit to data:
  - $P(n) = \frac{\gamma^K}{\Gamma(K)n!} \frac{\Gamma(n+K)}{(1+\gamma)^{n+K}}$
- $P(N, n)$  particle distribution for a cluster composed by  $N$  individual sources, usually taken to be poissonian.
- $W(N)$ : configuration, cluster size distribution:
  - $W(N) = \frac{\gamma}{\Gamma(K)} (\gamma N)^{K-1} \exp(-\gamma N)$



# Percolation of color sources

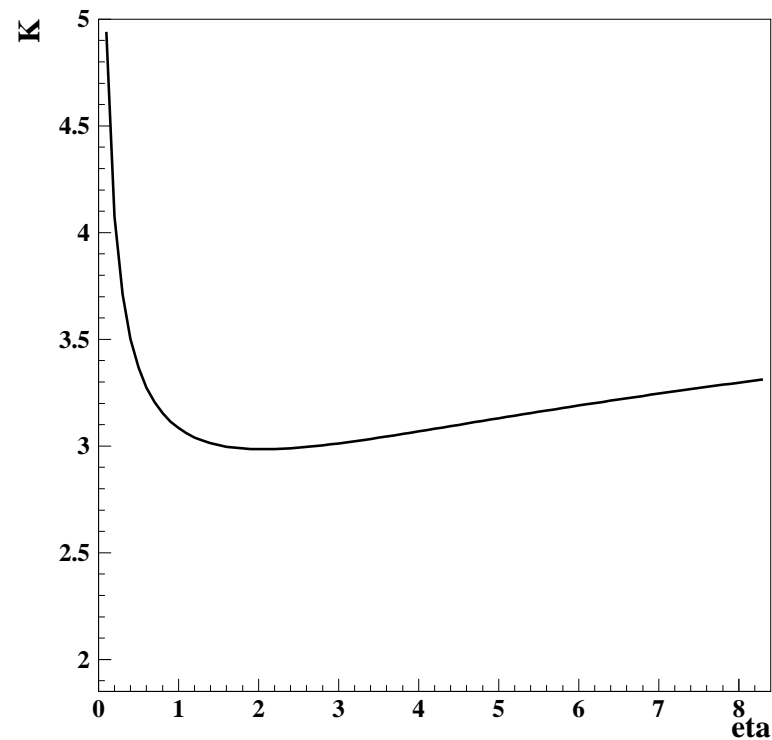


$$\frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle^2} = \frac{1}{K}$$
 ,K is a measure of the dispersion on the size of the emitting clusters

## Percolation of color sources

$$\frac{D(n)^2}{\langle n \rangle^2} = \frac{1}{K} + \frac{1}{\langle n \rangle}$$

$\eta = (\text{No.strings} * \text{stringsize}) / \text{totalsize}$



## Percolation of color sources

- $P_C(\nu) = \frac{\nu}{\langle \nu \rangle} P(\nu) \longrightarrow W_C(N) = \frac{N}{\langle N \rangle} W(N)$

- elementary collision  $\longrightarrow$  elementary cluster

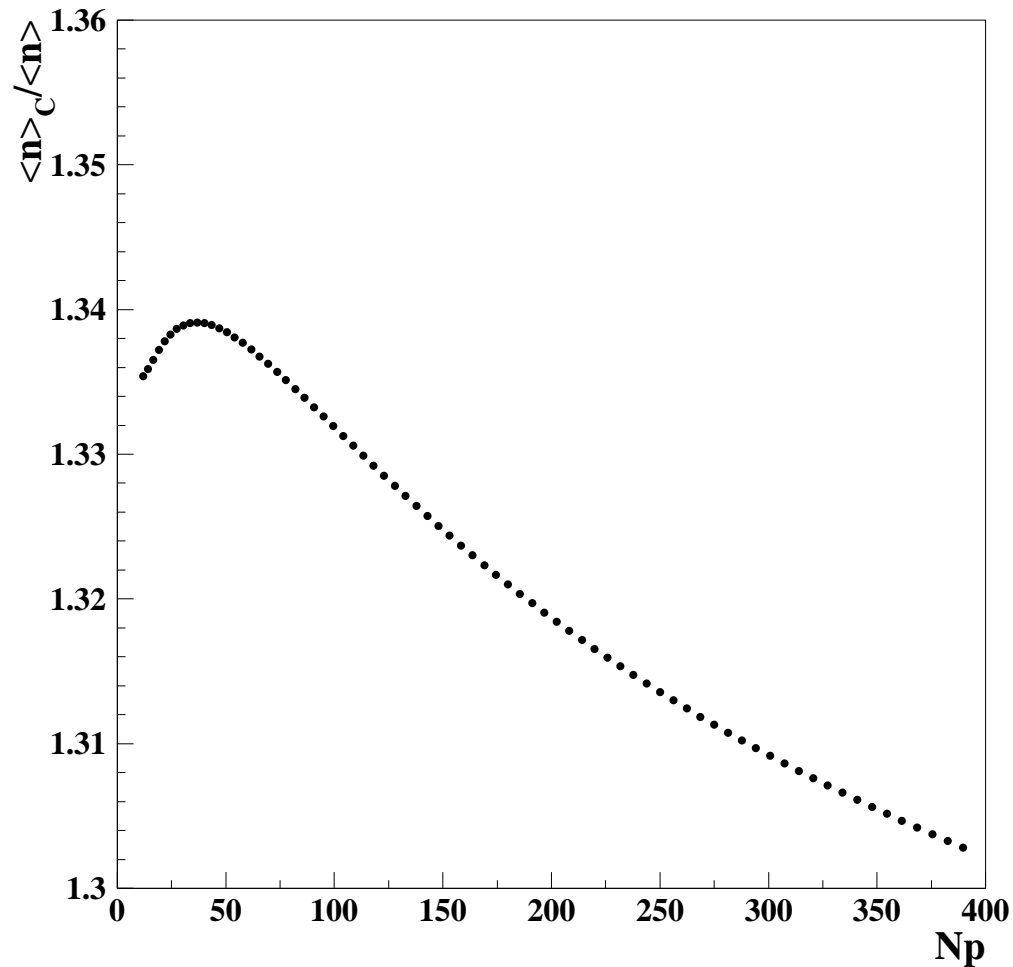
- $P_C(n) = \int_0^\infty dN \frac{N}{\langle N \rangle} W(N) P(N, n)$

- $P_C(n) = \frac{n+K}{\langle N \rangle} \frac{P(n)}{\gamma+1}$

- $\frac{\langle n \rangle_C}{\langle n \rangle} = \frac{1/K + 1/\langle n \rangle + K/\langle n \rangle + 1}{1 + K/\langle n \rangle}$

- for  $\langle n \rangle \gg K$ ,  $\frac{\langle n \rangle_C}{\langle n \rangle} = 1 + \frac{1}{K}$

$$\frac{\langle n \rangle_C}{\langle n \rangle} = 1 + \frac{1}{K} \text{ for Gold-Gold collisions at 200 GeV/A}$$



## Conclusions

- High  $p_T$  events at RHIC, heavy flavour production, particle production with no  $\delta y > \psi$ ...
- The same underlying mechanism

